

Laser-Raman spectrum of helical syndiotactic polypropylene

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A more intense and more detailed laser-Raman spectrum than previously reported for the two-fold helix conformation of syndiotactic polypropylene is presented. Good agreement is found between the observed Raman shifts and those predicted by other authors; hence assignment to symmetry species of much of the experimental data is possible.

INTRODUCTION

In their review Youngman and Boor¹ emphasize the need for a complete characterization study of polypropylene, in order to understand and interrelate polymer structure property relationships. Since all steric arrangements can be prepared, from syndiotactic through atactic to isotactic, the study of polypropylene has special relevance to an elucidation of Ziegler-type catalysis and polymer behaviour¹.

The vibrational spectrum of isotactic polypropylene has been the subject of many papers. These include normal coordinate calculations and their comparison with experimental infra-red data²⁻⁸, infra-red spectra⁸⁻¹¹ and Raman data¹²⁻¹⁵. Vibrational frequencies for both crystalline forms of the syndiotactic polymer (SPP) have been determined from force-field calculations derived from a series of saturated hydrocarbons⁹ by Schachtsneider and Snyder¹⁶ and transferred from isotactic polypropylene³ by Miyazawa¹⁷ and Miyazawa and Ideguchi¹⁸. These calculated frequencies have been compared with the experimental infra-red data of Peraldo and Cambini¹⁹ and Tadokoro *et al.*²⁰. Koenig *et al.*²¹ have also compared the infra-red spectra of the helical forms

of both isotactic and syndiotactic polypropylene with that of the atactic material.

Previous reports of experimental Raman data for helical SPP are limited to those of Zerbi and Hendra²² for a highly crystalline powder sample and a brief reference by Zerbi²³ to a 'regularity band' in the skeletal vibration region near $850 \Delta\text{cm}^{-1}$. Unfortunately the observation of Raman signals by Zerbi and Hendra²² was severely hampered by fluorescence, and consequently only the most intense Raman bands, ascribed to A modes, were evident in their spectrum. In addition their recorded spectrum exhibited a strong emission for which no reasonable explanation was offered other than that fluorescence was responsible. In this paper we present a Raman spectrum (*Figure 1*) of syndiotactic polypropylene in the two-fold helix conformation which is a substantial improvement over that previously published, being more intense and also free of fluorescence background. The data recorded allow a more complete observation of Raman signals in the region $180-4000 \Delta\text{cm}^{-1}$ and a better comparison with the predicted vibrational frequencies^{16,17}. The good agreement found between observed and predicted frequencies permits assignment to symmetry species of much of the experimental data.

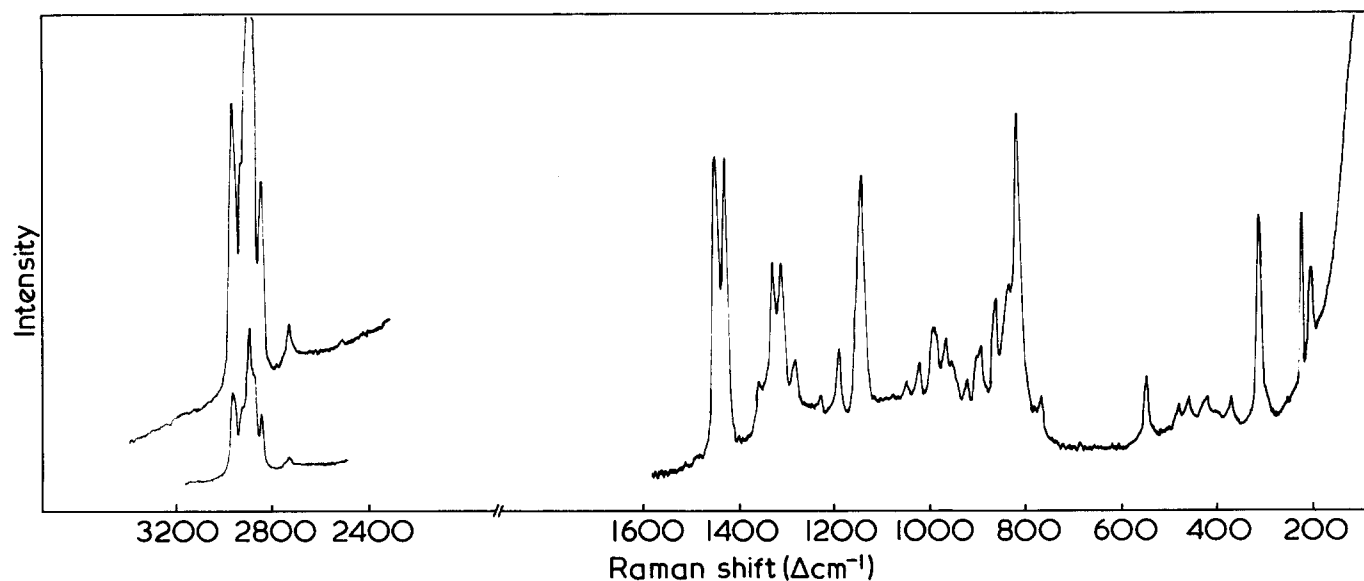


Figure 1 Laser-Raman spectrum for a crystalline powder sample of helical syndiotactic polypropylene

EXPERIMENTAL

The spectrum shown in Figure 1 was recorded from a crystalline powder sample using 488.0 nm excitation from an argon ion Lexcel model 75 laser (~80 mW power at the sample) on a Cary 83 laser-Raman spectrometer. The slit width was set at 5 cm⁻¹, the pen period at 10 sec and the scan speed at 0.2 Δcm⁻¹/sec. All the bands excepting that at ~220 Δcm⁻¹ (plasma emission) arose from Raman emission by the sample. The poor stray-light discrimination properties of the small 'bench-top' Cary 83 Raman spectrometer prevented the observation of any signals at frequencies below 180 Δcm⁻¹.

The polymer was characterized by ¹³C n.m.r. and X-ray spectroscopic techniques. The triad sequences of ¹³C n.m.r. spectra of polypropylene are well characterized²⁴ and for this sample only rr- and mr-centred²⁵ units were detected from the spectrum. The former triads represented about 90% of the observed resonances. Since there was no evidence for mm-centred units the latter triads must be present as rmr sequences. It may therefore be considered to be an SPP sample of high purity. The X-ray powder photograph was characteristic²⁶ of the two-fold helix conformation of SPP²⁷.

DISCUSSION

Syndiotactic polypropylene can be prepared in two crystalline modifications²⁸: the more stable helical conformation, and a planar zig-zag form. The helical conformation, the subject of this paper, has a polymer chain with four monomeric units per translational repeat distance and three mutually perpendicular two-fold axes, the helix axis and two-fold axes passing through the CH₂ groups, and has a conformation described¹⁶ as TGGTTG'G'T. The line group of this model is isomorphous with the point group D₂ and there are 26 spectroscopically active vibrations^{16,17} for each of the species A, B₁, B₂ and B₃. The A modes are Raman active only. Schachtsneider and Snyder¹⁶ define the B₁, B₂ and B₃ species as those modes with transition moments parallel to the two-fold axis passing through the external CH₂, the two-fold axis passing through the internal CH₂ and the chain axis respectively. Where quoted, the species as defined by Miyazawa¹⁷ have been transformed to agree with this notation. Although there is reasonable agreement between the calculated frequencies of Schachtsneider and Snyder¹⁶ and Miyazawa¹⁷, the values quoted by the former authors are used to compare to the Raman shifts observed in this study. The more complete assignment and good agreement obtained in the infra-red by a comparison of their predictions with the studies of Peraldo and Cambini¹⁹ on bi-dimensionally oriented specimens leads us to favour their assignments.

VIBRATIONAL ASSIGNMENTS

In the absence of any other data (anisotropic scattering, temperature dependence)¹², and because of the many accidental degeneracies arising from the weak intramolecular coupling²², it is not possible to assign unambiguously the observed Raman bands to vibrations of A, B₁, B₂ and B₃ species. However, the good agreement found by Schachtsneider and Snyder¹⁶ between observation and prediction for the infra-red spectrum, and the frequency coincidence or near coincidence of many of the Raman bands seen in this study with the infra-red absorption bands deter-

Table 1 Observed and calculated frequencies for the D₂ form of SPP

ν _{Calc} (cm ⁻¹)			ν _{Obs} (cm ⁻¹)	
Ref 17	Ref 16		Raman	
		Infra-red Ref 16*	This study	Ref 22
1463	1463	A	{1466 s	{1463 w
1458	1462	A		
1438	1454	A		
1435	1452	A	{1445 s	{1445 m-w
1462	1465	B ₁		
1458	1462	B ₁	{1465 1460 1432 1463	
1437	1453	B ₁		
1462	1464	B ₂		
1461	1461	B ₂		
1437	1454	B ₂		
1461	1464	B ₃	{1460 1455	
1458	1462	B ₃		
1387	1380	B ₁	1379	
1384	1374	B ₂	1378?	
1384	1378	B ₃	{1377 1374	{1374 w
1355	1373	B ₃		
1363	1372	B ₂	1373	
1387	1372	A	1370 sh,w?	
1366	1357	B ₁		
1368	1353	A	1345 m	1346 vs
1329	1352	B ₁	1332	
1349	1352	B ₃	1346	
1338	1339	A	1327 m	1331 vs
1345	1331	B ₂	1360	
1287	1301	B ₂	1311	
1300	1300	B ₃	1293	1297 w
1289	1275	B ₁	1287	
1263	1265	B ₃	1264	
1211	1259	A		
1190	1223	B ₁	1242	1244 vw
1195	1194	B ₂	1202	
1190	1191	A		1204 w-m
1147	1174	B ₃	1167	1202 vw
1139	1168	A		1157 s
1147	1167	B ₁	1152	1158 m
1148	1159	B ₂	1153	
1115	1115	B ₂		
1113	1113	B ₁	1088, 1083	
1062	1096	B ₃		
1057	1068	B ₃	1060	1061 vw
1033	1044	B ₂	1035	{1034 w
1112	1041	A		1040 vw
1007	996	A		1002 m
979	980	B ₁	1006	1000 vw
948	974	B ₃	977	
939	971	B ₁	976? a	{978 w-m
			963 a†	966 w
929	931	B ₂	935	933 vw
926	918	A		913 sh,w
901	903	B ₁	906	
904	902	B ₃	901	{903 w
883	878	B ₂	~870	872 m
873	867	B ₃	867	
			845 †	844 sh,w
857	839	A		828 s
820	829	B ₃		829 m-s
826	825	B ₂	812	
824	786	B ₁	776	776 w
529	542	A		552 w-m
545	537	B ₁	535	
484	476	B ₂	483	486 w
469	436	B ₃	468	463

v = very; w = weak; m = medium; s = strong; sh = shoulder; a = amorphous.

* The frequencies listed in this column are those given by Schachtsneider and Snyder¹⁶ from data supplied by Natta, Peraldo and Cambini prior to their own publication¹⁹. Slight discrepancies, nominally not more than ±1 Δcm⁻¹, exist between these two sets of figures. The most notable differences are for two B₂ bands quoted as 1153 and 1035 cm⁻¹ in reference 16 and as 1160 and 1032 cm⁻¹ respectively in reference 19. Potential energy distributions and descriptions of the vibrations are thoroughly covered in reference 16.

† Data from reference 19

Table 2 Observed and calculated frequencies for the D_2 form of SPP, annotations as for Table 1

$\nu_{\text{Calc}} (\text{cm}^{-1})$			$\nu_{\text{Obs}} (\text{cm}^{-1})$	
Ref 17	Ref 16	Infra-red Ref 16*	Raman	
			This study	Ref 22
2963,2962	2962,2962	A	2962 s	2960 w
2963,2962	2962,2962	B ₁	2922 sh,w	
2963,2962	2962,2962	B ₂		
2963,2962	2962,2962	B ₃		
2903	2929	B ₁	2887 vs	2904 m
2904	2929	B ₂		
2906	2930	B ₃	2873 sh, w-m	2892-2872 m
2902	2928	B ₃		
2926	2906	A		
2927-2926	2904,2904	B ₁ ,B ₂	2840 m	2848 w-m
	2903	B ₃		
2873	2882	A	2725 w	
2872-2873	2882	B ₁		
	2882	B ₂		
	2882	B ₃		
2849	2856	A		
2844	2855	A		
2847	2855	B ₁		
2847	2855	B ₂		

mined by Peraldo and Cambini¹⁹, suggest origins for much of the Raman data. (Schachtsneider and Snyder¹⁶ were able to assign nearly all of the observed absorption bands¹⁹ and with the exclusion of the C-H stretching frequencies and the calculated frequency of the symmetric <CCC bending, which was obviously out of line¹⁶, their average error between the two sets of data was 0.85%.) We have also adhered to the criterion of Zerbi and Hendra²² that the totally symmetric A modes are most likely to give rise to the most intense Raman bands. Bands observed only in the Raman spectrum are also likely to be A modes. This makes possible assignments of the majority of the observed Raman bands.

THE REGION 1500-450 Δcm^{-1}

Origins or major contributors to the intensity of the Raman bands observed for the region 1500-450 Δcm^{-1} are given in Table 1. At the high frequency end of this range, distinct splitting of each vibration into four easily separable species is poor; however, as expected²², most of the observed Raman emission intensity between 1100 and 1500 Δcm^{-1} seems to be attributable to A modes. Near-neighbour B species may well contribute to the asymmetry seen on the high frequency side of the 1157 Δcm^{-1} band and implied in the complex emission centred at about 1335 Δcm^{-1} . It is difficult, however, to reconcile any of the observations with origination from the A mode prediction at 1259 cm^{-1} , the nearest (very weak) observation at 1244 Δcm^{-1} seeming to favour assignment to the species B₁. The weak feature at 966 Δcm^{-1} does not appear to correspond to a normal mode of vibration for helical SPP. It is best assigned as arising from zig-zag regions of the SPP chain in the amorphous state¹⁹. A similar weak feature was observed in the infra-red spectrum¹⁹ and strong bands near 963 Δcm^{-1} were seen in the infra-red spectra of the zig-zag conformation of SPP and for SPP in solution and in the molten state. The weak shoulder at 844 Δcm^{-1} corresponds with the weak infra-red absorption at 845 cm^{-1} and is associated with non-helical conformations^{19,23}.

THE REGION 450-180 Δcm^{-1}

Below 450 Δcm^{-1} four bands were observed: 424 (weak), 376 (weak), 315 (strong) and 202 (medium). The intense band at 315 Δcm^{-1} must be assigned to the A mode described as a mixture of <C-C-C bend and C-C skeletal stretch^{16,22}, the predicted frequency of which is coincident with the observation. The medium intensity band at 202 Δcm^{-1} , which originates from CH₃ torsion, is best attributed to the A mode (199 cm^{-1}); the corresponding B₁ (196), B₂ (202) and B₃ (200) species possibly contribute to its intensity. Schachtsneider and Snyder¹⁶ found that their calculated frequency of the symmetric <CCC bending of the C(C₃) group was obviously out of line and the B₃ mode predicted at 436 cm^{-1} was found in the infra-red at 468 cm^{-1} . (Miyazawa¹⁷ calculated this mode to occur at 469 cm^{-1} .) The contribution of this symmetric bend to the potential energy distribution¹⁶ was 46%; similar errors might then be expected for the B₂ mode (37%) predicted at 433 cm^{-1} , the B₁ mode (40%) at 403 cm^{-1} and the A mode (27%) at 347 cm^{-1} . Both the above B species might therefore be considered as candidates for the weak feature at 424 Δcm^{-1} with the A mode possibly giving rise to the observation at 376 Δcm^{-1} . (The only other predictions in this region were B₃ modes at 410 cm^{-1} and 358 cm^{-1} .) Amorphous or intermediate conformations cannot be precluded as alternative origins for these two bands. Zerbi and Hendra²² also observed low frequency bands near 315 Δcm^{-1} and 200 Δcm^{-1} , at 314 Δcm^{-1} (medium strong) and 204 Δcm^{-1} (weak) respectively and an additional weak feature at 172 Δcm^{-1} (A-mode predicted¹⁶ at 176 cm^{-1}).

THE REGION 3000-2700 Δcm^{-1}

The data for this region are given in Table 2. The CH-stretching region appears again to be dominated by A modes with the 2922 Δcm^{-1} observed Raman shift originating from one or several coincident B species. The 2725 Δcm^{-1}

observation probably arises as an overtone or combination of the intense bands observed in both the infra-red and Raman spectra near 1450 cm^{-1} and 1350 cm^{-1} .

CONCLUSION

The laser Raman spectrum for the two-fold helix conformation of syndiotactic polypropylene recorded shows many more vibrational bands than previously reported. The good correlation found between observation and prediction verifies further the force field of Schachtsneider and Snyder¹⁶ and proposed polymer structure.

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